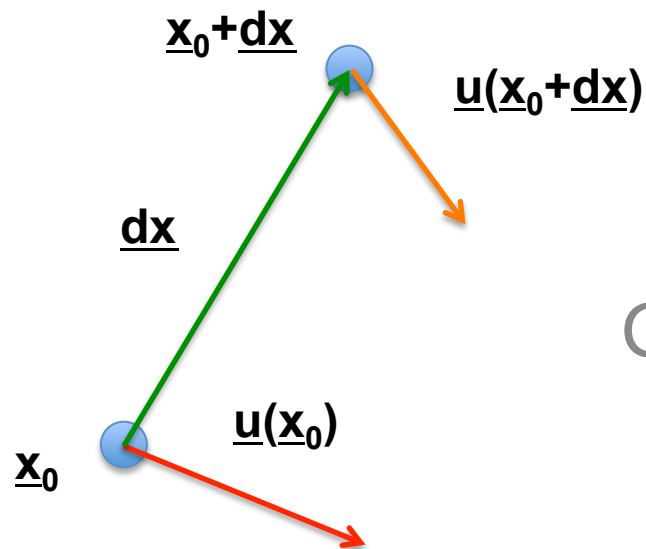


# Lecture 12: Coordinates, Transformations, vectors, displacements, strains

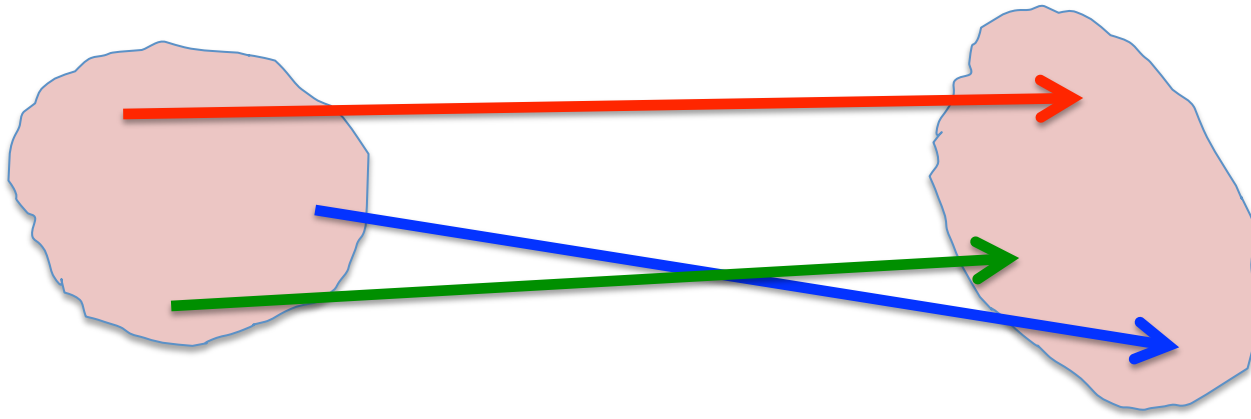


GEOS 655 Tectonic Geodesy  
Jeff Freymueller

# Outline

- Coordinates and Transformations
  - The mathematical basis for the problem
  - In particular, how to do rotations
- Motions in general
  - displacement = rigid body motion + deformation
  - displacement = translation + rotation + deformation
- Deformation = strain

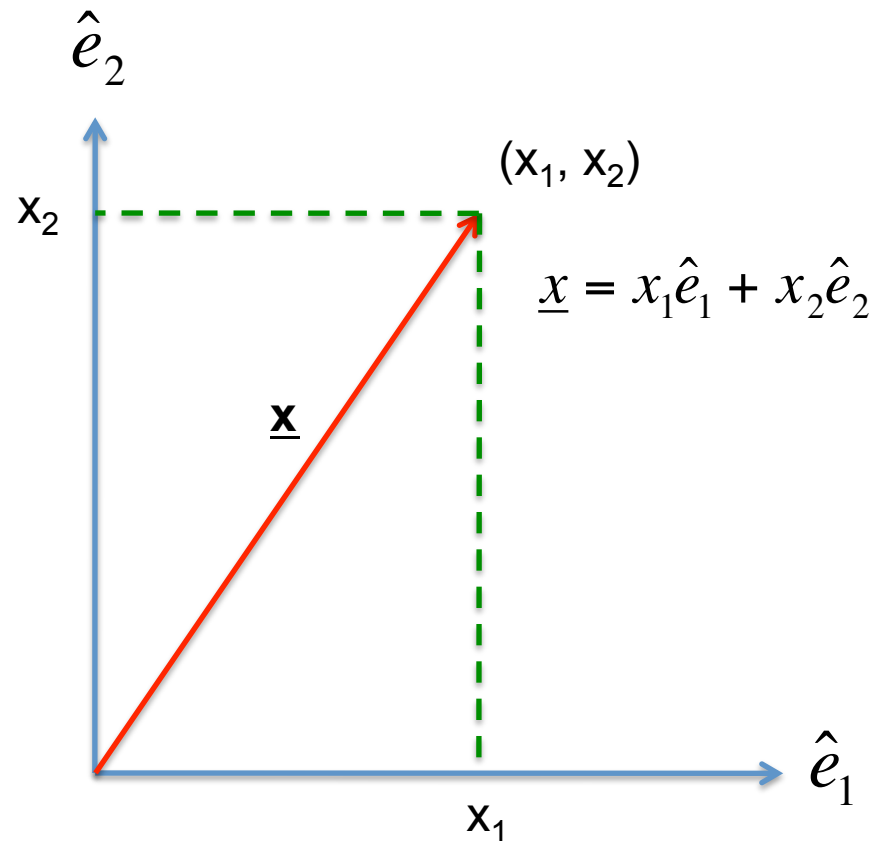
# Motion in General



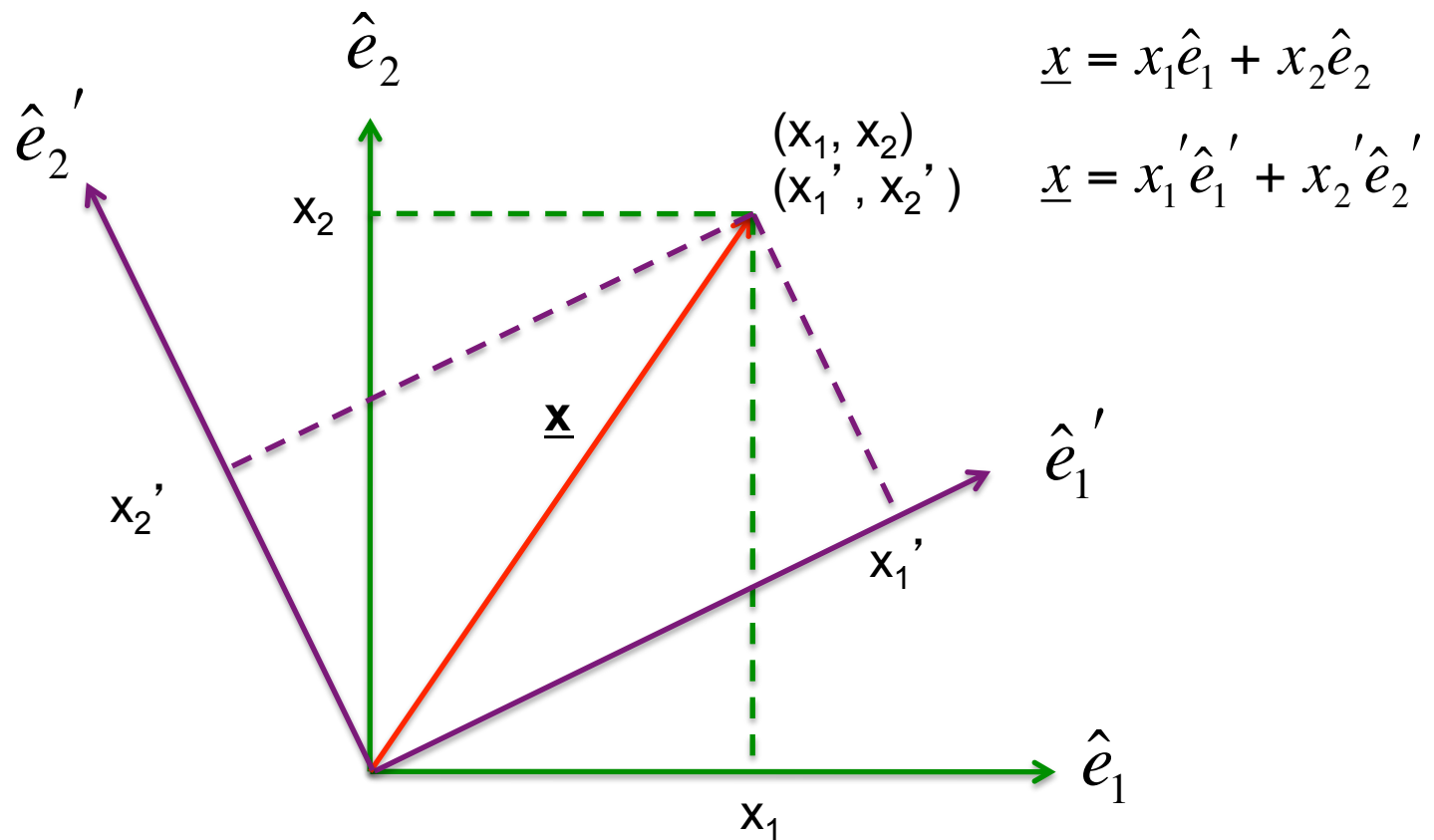
- displacement = rigid body motion + deformation
- displacement = translation + rotation + deformation

# Coordinates and Vectors (in 2D)

- Given a point and coordinate axes, we can define a pair of **coordinates** that locate the point.
  - Vector may be a physical thing
  - Coordinates represent the vector
  - Imagine a vector that points from the center of the earth to a point on the surface. If we change coordinate systems, the vector still points to the same place, but the coordinates or vector components will be different.



# Same Vector, Different Coordinates



Mathematically, we often blur the distinction between the vector and its representation.

# Many Different Vector Notations

$$\underline{x} = (x_1, x_2, x_3) = \begin{bmatrix} x_1 & x_2 & x_3 \end{bmatrix} = \begin{bmatrix} x_1 \\ x_2 \\ x_3 \end{bmatrix} \quad \begin{array}{l} \text{row vector} \end{array} \quad \text{column vector}$$

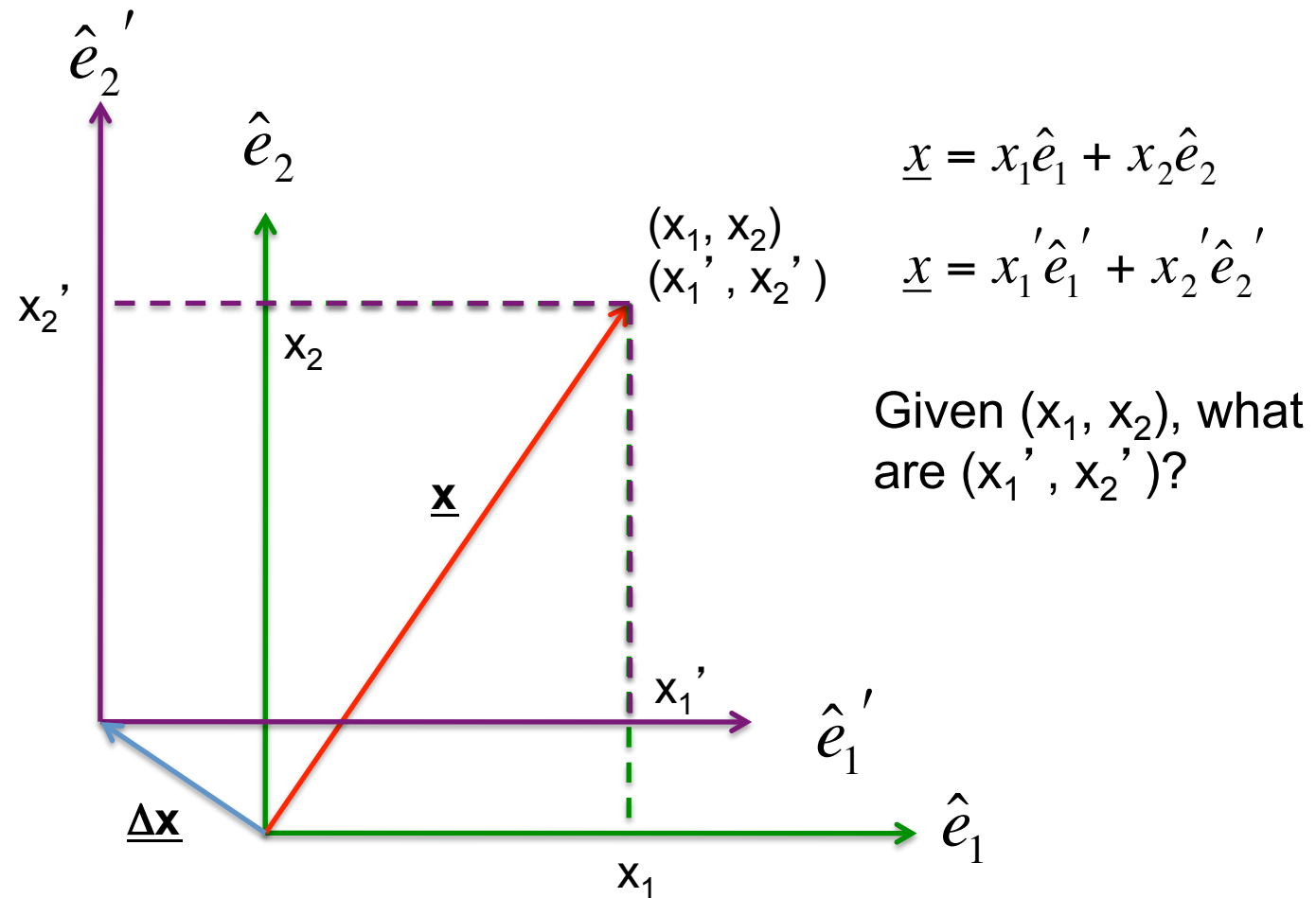
- This leads to a variety of notations for vector operations, for example the dot product:

$$\begin{aligned} \underline{x} \cdot \underline{y} &= x_1 y_1 + x_2 y_2 + x_3 y_3 \\ &= \sum_{i=1}^3 x_i y_i \\ &= x_i y_i \quad \leftarrow \text{"Einstein convention": summation over repeated indices} \\ &= \underline{x}^T \underline{y} \\ &= \begin{bmatrix} x_1 & x_2 & x_3 \end{bmatrix} \cdot \begin{bmatrix} x_1 \\ x_2 \\ x_3 \end{bmatrix} \end{aligned}$$

# Transformations

- Mathematical ***transformations*** map something from one system (or set) to another
  - For example, transform the coordinates of a vector from one coordinate system to another
  - A general form:  $\underline{y} = T\underline{x}$
- Linear transformations are part of the basis of ***linear algebra***.
  - A mathematician would say that a transformation is a mapping from one “space” to another.
- We can represent linear transformations in a variety of ways, but the two most common are a matrix representation or using index notation

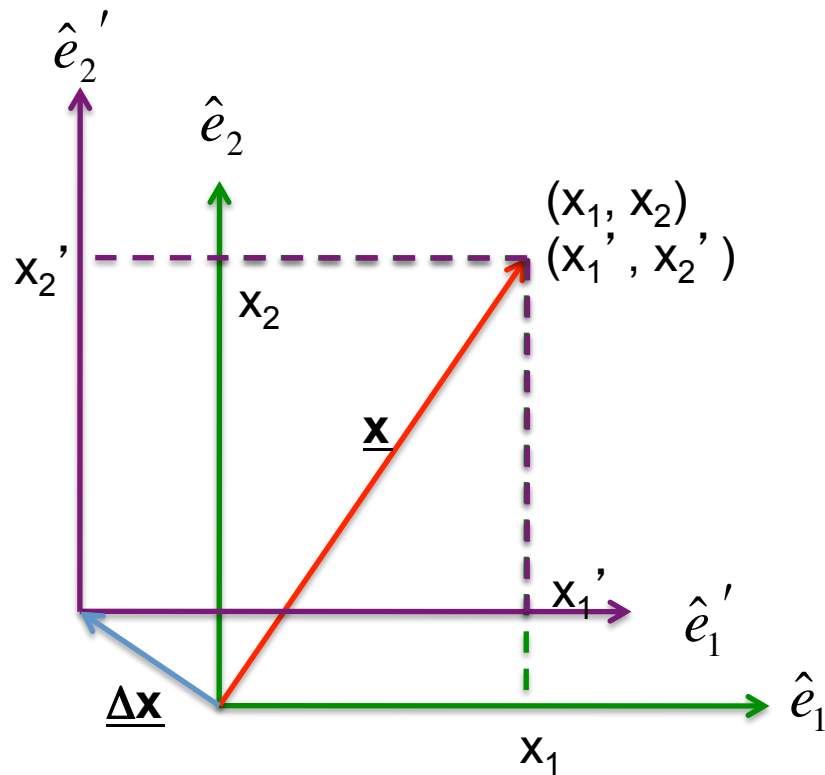
# Translational Transformations



In this case the unit vectors describing the coordinate system are the same, only the origin of the system is different.



# Translational Transformations



$$\underline{x} = x_1 \hat{e}_1 + x_2 \hat{e}_2$$

$$\underline{x} = x_1' \hat{e}_1' + x_2' \hat{e}_2'$$

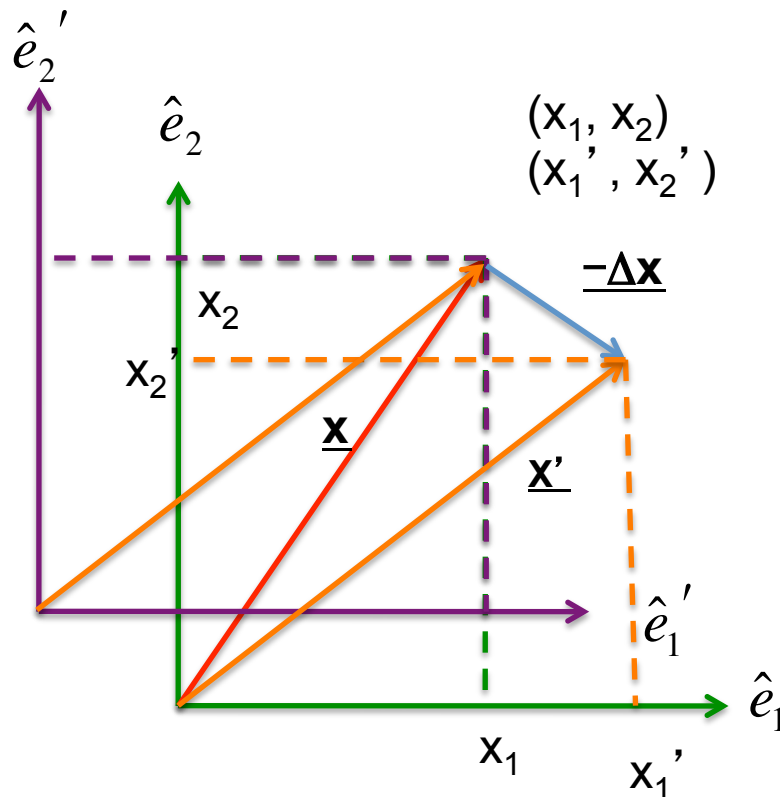
Given  $(x_1, x_2)$ , what are  $(x_1', x_2')$ ?

This is pretty easy:

$$x_1' = x_1 - \Delta x_1$$

$$x_2' = x_2 - \Delta x_2$$

# Translational Transformations



$$\underline{x} = x_1 \hat{e}_1 + x_2 \hat{e}_2$$

$$\underline{x}' = x'_1 \hat{e}'_1 + x'_2 \hat{e}'_2$$

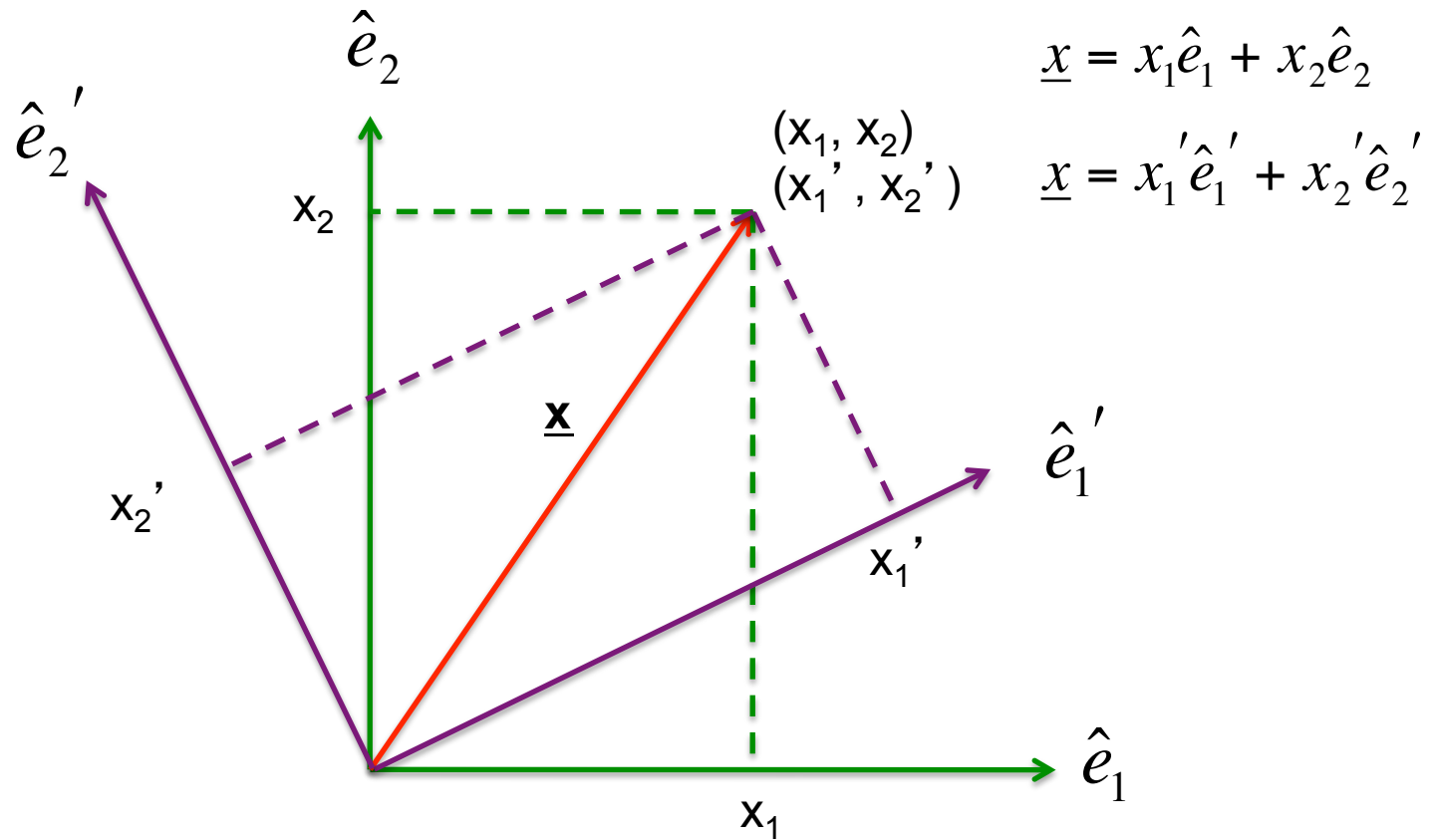
$$x'_1 = x_1 - \Delta x_1$$

$$x'_2 = x_2 - \Delta x_2$$

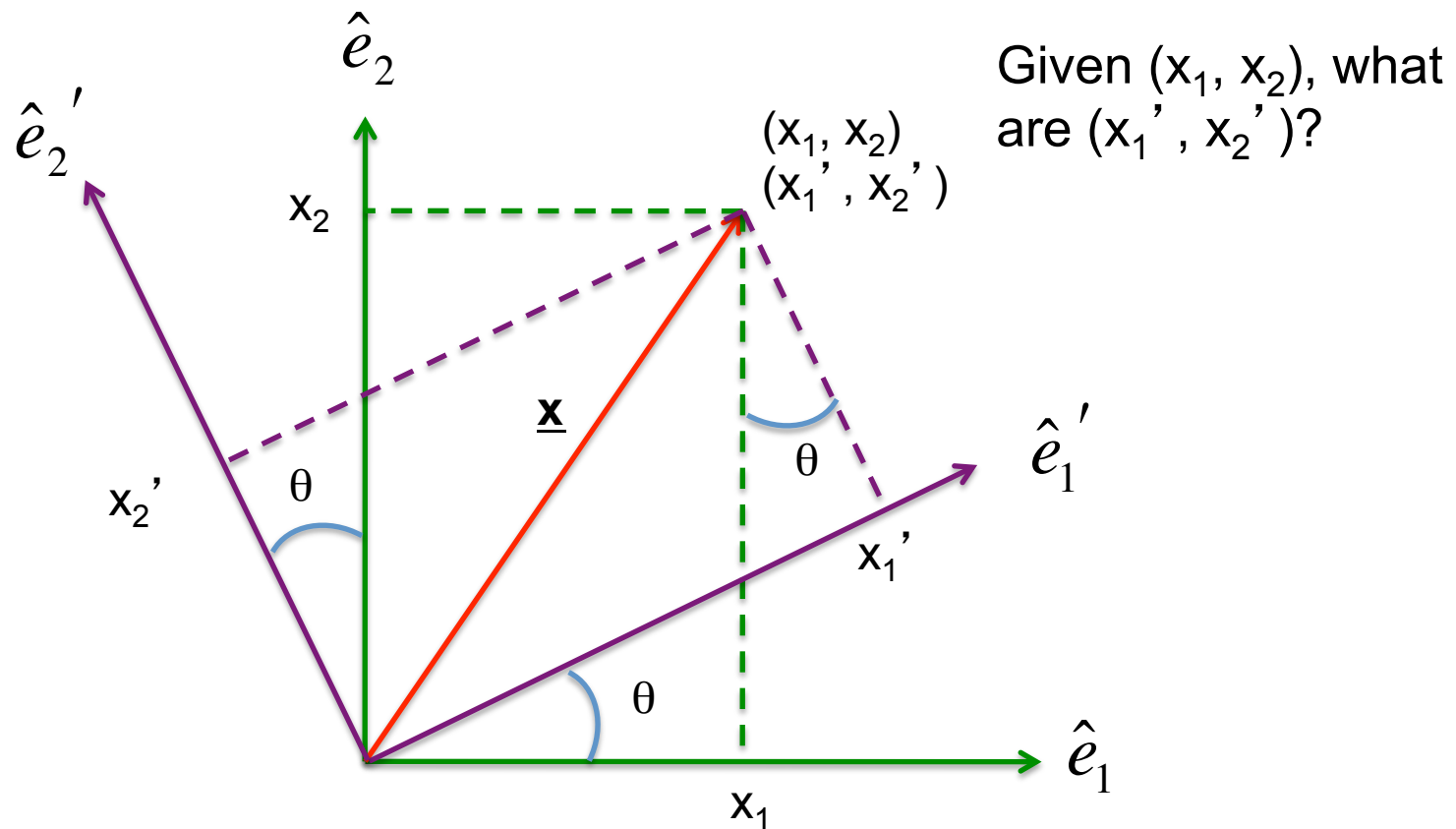
This vector  $\underline{x}'$  is also given by  $\underline{x}' = \underline{x} - \underline{\Delta x}$

Translating the coordinate system by  $\underline{\Delta x}$  is like subtracting  $\underline{\Delta x}$  from the vector.

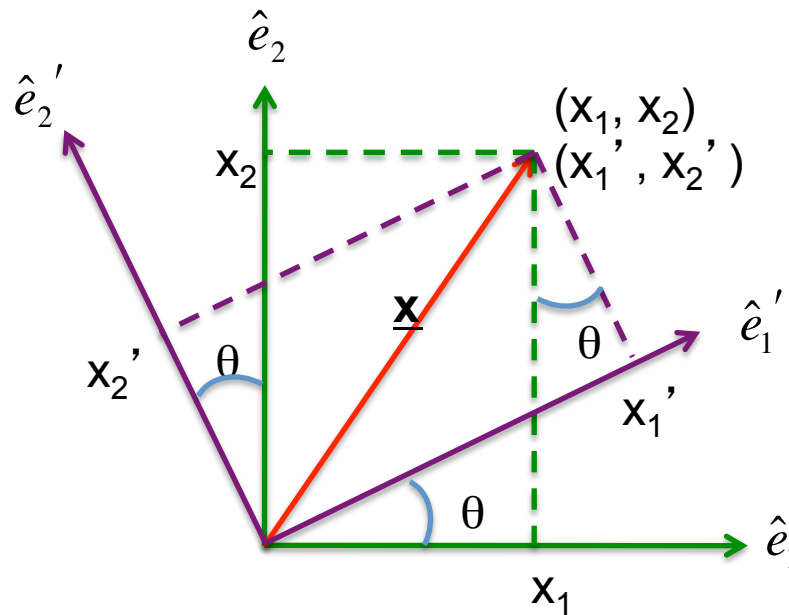
# Rotational Transformations



Given two systems, how are the components of a vector related?



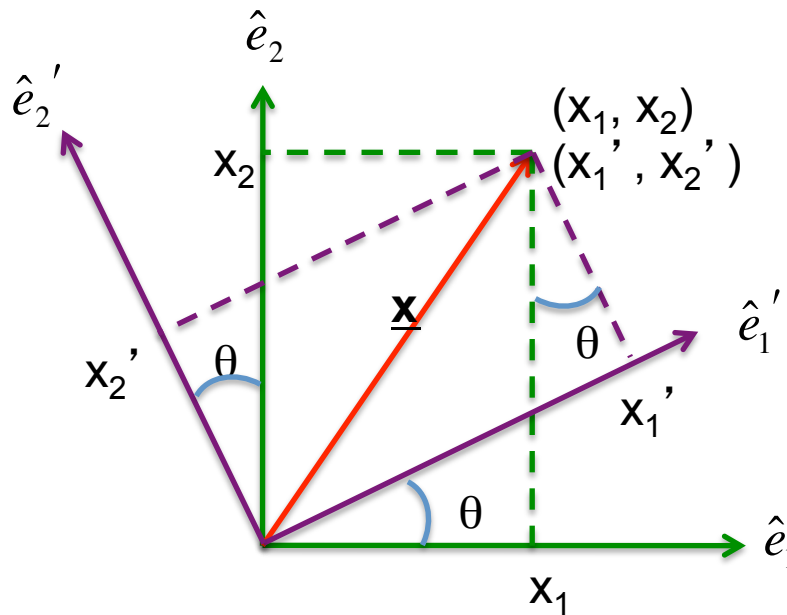
Given two systems, how are the components of a vector related?



Given  $(x_1, x_2)$ , what are  $(x_1', x_2')$ ?

$$\begin{aligned} x_1 &= x_1' \cos \theta - x_2' \sin \theta & x_1' &= x_1 \cos \theta + x_2 \sin \theta \\ x_2 &= x_1' \sin \theta + x_2' \cos \theta & x_2' &= -x_1 \sin \theta + x_2 \cos \theta \end{aligned}$$

Given two systems, how are the components of a vector related?



Given  $(x_1, x_2)$ , what are  $(x_1', x_2')$ ?

$$x_i = \beta_{ij} x_j' \quad \text{implied summation}$$

$$(\beta_{ij}) = \begin{bmatrix} \beta_{11} & \beta_{12} \\ \beta_{21} & \beta_{22} \end{bmatrix} = \begin{bmatrix} \cos \theta & -\sin \theta \\ \sin \theta & \cos \theta \end{bmatrix}$$

# Transformation and Inverse

- The matrix of the betas is the ***rotation matrix***, which maps coordinates in the primed system into coordinates in the unprimed system:

$$x_i = \beta_{ij} x_j'$$

- The ***inverse transformation*** (opposite rotation) goes the other way, from the unprimed to the primed system

$$x_i' = \beta_{ji} x_j = \beta_{ki} x_k$$

$$(\beta_{ji}) = (\beta_{ij})^T = \begin{bmatrix} \beta_{11} & \beta_{21} \\ \beta_{12} & \beta_{22} \end{bmatrix} = \begin{bmatrix} \cos \theta & \sin \theta \\ -\sin \theta & \cos \theta \end{bmatrix}$$


# Transformation and Inverse

- Combining the two equations:

$$x_i = \beta_{ji} x_j' = \beta_{ji} (\beta_{jk} x_k) = (\beta_{ji} \cdot \beta_{jk}) x_k$$

$$(\beta_{ji} \cdot \beta_{jk}) = \delta_{ik}$$

$$\begin{bmatrix} x_1 \\ x_2 \end{bmatrix} = \begin{bmatrix} \beta_{11} & \beta_{21} \\ \beta_{12} & \beta_{22} \end{bmatrix} \cdot \begin{bmatrix} x_1' \\ x_2' \end{bmatrix}$$

*matrix form*   $\begin{bmatrix} x_1 \\ x_2 \end{bmatrix} = \begin{bmatrix} \beta_{11} & \beta_{21} \\ \beta_{12} & \beta_{22} \end{bmatrix} \cdot \left( \begin{bmatrix} \beta_{11} & \beta_{12} \\ \beta_{21} & \beta_{22} \end{bmatrix} \cdot \begin{bmatrix} x_1 \\ x_2 \end{bmatrix} \right)$

$$\begin{bmatrix} \beta_{11} & \beta_{21} \\ \beta_{12} & \beta_{22} \end{bmatrix} \cdot \begin{bmatrix} \beta_{11} & \beta_{12} \\ \beta_{21} & \beta_{22} \end{bmatrix} = I$$



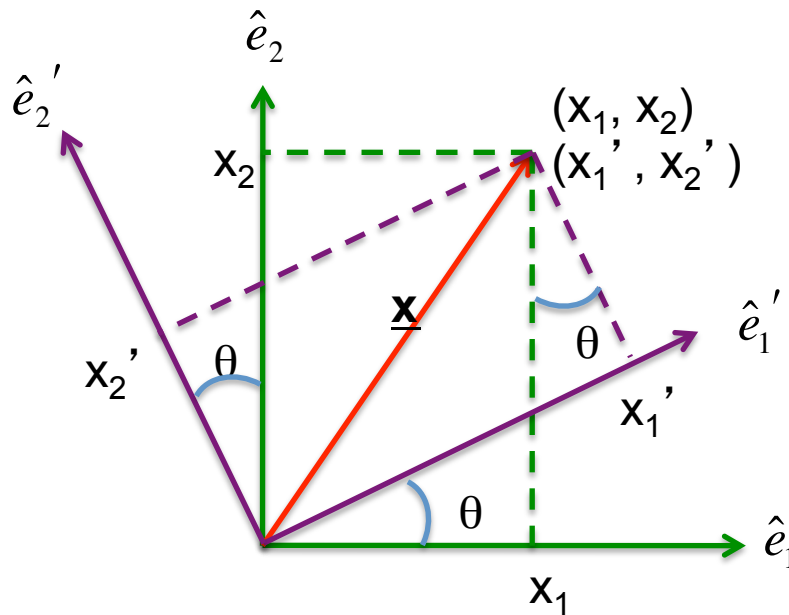
# Transformation and Inverse

- The transpose of the rotation matrix is its inverse.

$$\begin{bmatrix} \beta_{11} & \beta_{12} \\ \beta_{21} & \beta_{22} \end{bmatrix}^T = \begin{bmatrix} \beta_{11} & \beta_{12} \\ \beta_{21} & \beta_{22} \end{bmatrix}^{-1}$$

- This kind of matrix is called an ***orthogonal matrix***.
- Rotational transformations preserve length
  - The length of the vector is the same in both coordinate systems.
  - Other kinds of linear transformations can be done that do not preserve length.

# Geometric Relationships



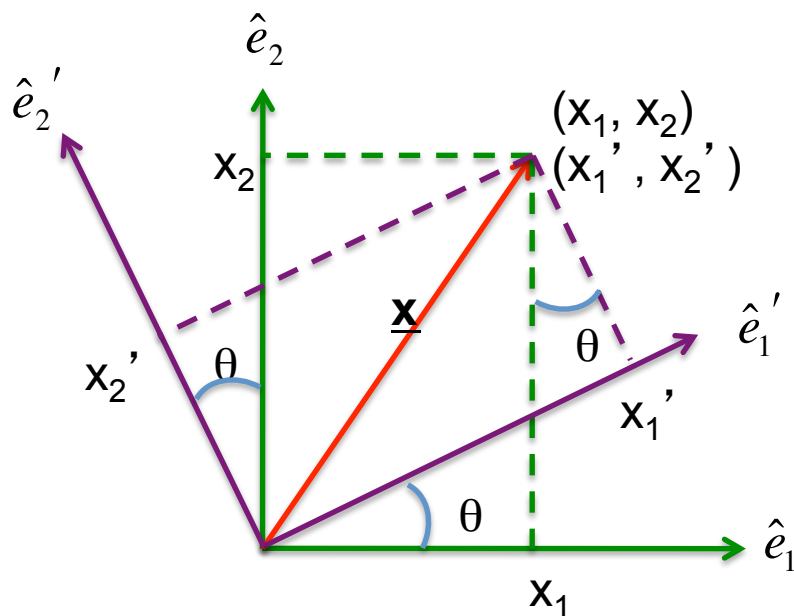
$$\beta_{11} = \cos \theta = \hat{e}_1 \cdot \hat{e}_1'$$

$$\beta_{ij} = \hat{e}_i \cdot \hat{e}_j'$$

- The betas have an easy geometric interpretation
  - Sometimes called **direction cosines**

$$(\beta_{ij}) = \begin{bmatrix} \beta_{11} & \beta_{12} \\ \beta_{21} & \beta_{22} \end{bmatrix} = \begin{bmatrix} \cos \theta & -\sin \theta \\ \sin \theta & \cos \theta \end{bmatrix}$$

# Geometric Relationships



$$\underline{x} = x_j \hat{e}_j = x_j' \hat{e}_j'$$

- Take the dot product of both sides with  $\hat{e}_i$

$$(x_j \hat{e}_j) \cdot \hat{e}_i = (x_j' \hat{e}_j') \cdot \hat{e}_i$$

$$x_j (\hat{e}_j \cdot \hat{e}_i) = x_j' (\hat{e}_j' \cdot \hat{e}_i)$$

$$\delta_{ij} x_j = (\hat{e}_i \cdot \hat{e}_j') x_j'$$

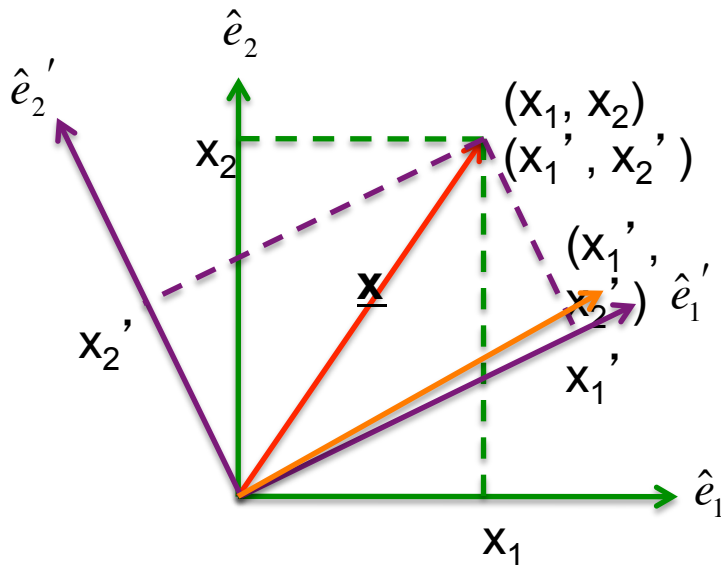
$$x_i = (\hat{e}_i \cdot \hat{e}_j') x_j' = \beta_{ij} x_j'$$

# Rotating Coordinate System vs. Rotating the Vector

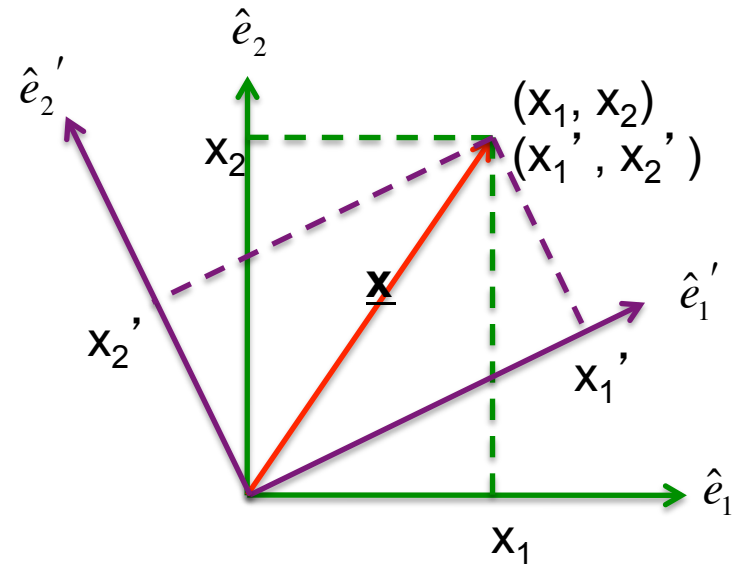
Rotating a vector is the same as rotating the coordinate system in the opposite direction

$$\underline{x} = x_1 \hat{e}_1 + x_2 \hat{e}_2$$

$$\underline{x} = x_1' \hat{e}_1' + x_2' \hat{e}_2'$$

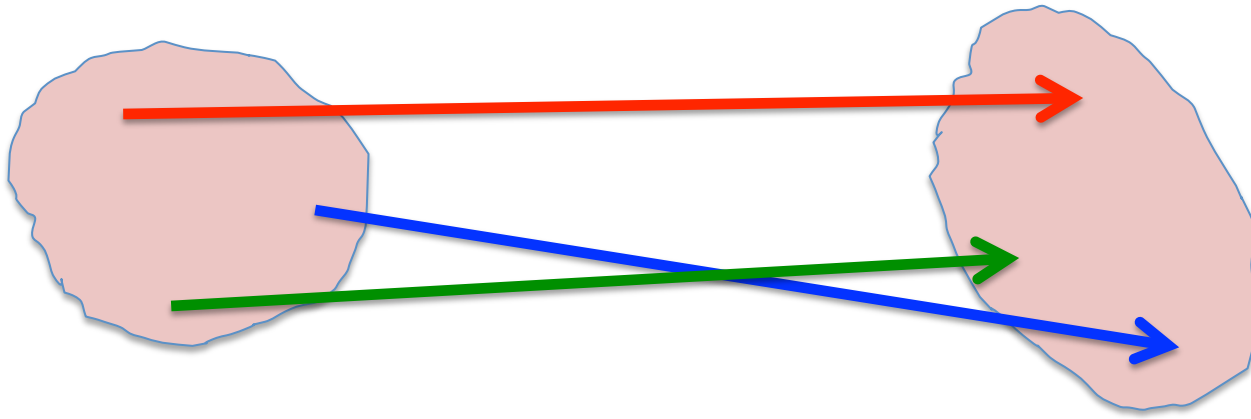


Clockwise rotation of vector



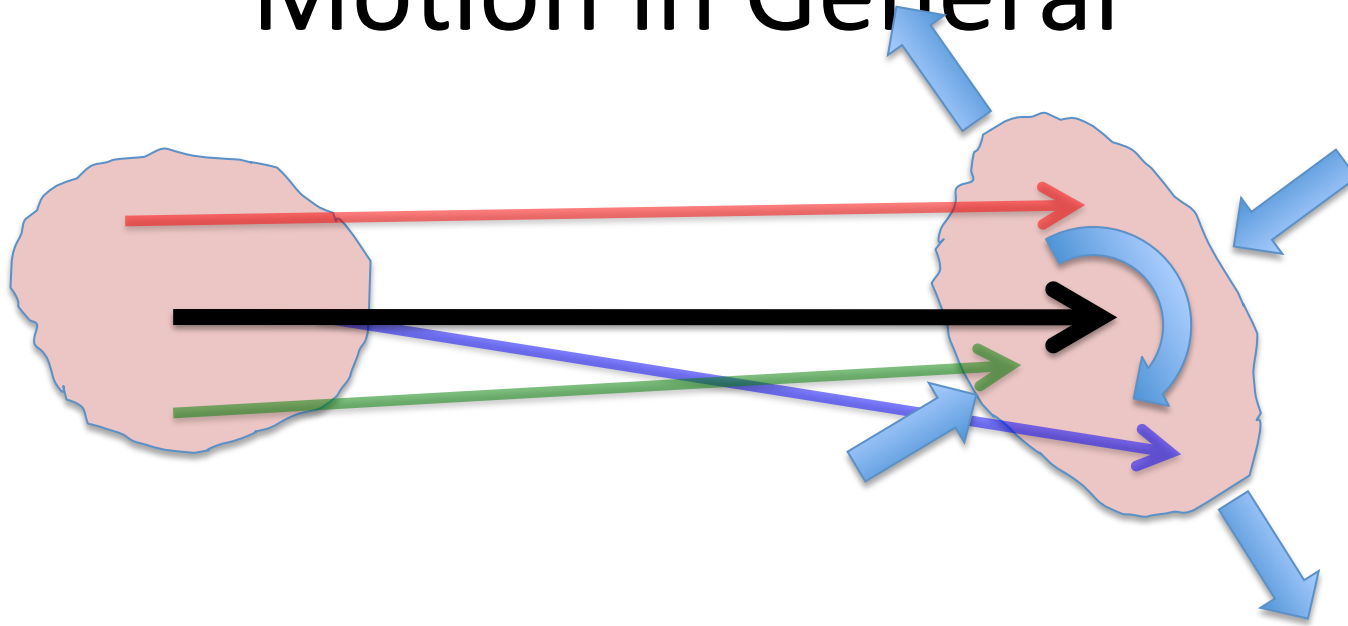
Counter-clockwise rotation of coordinate system

# Motion in General



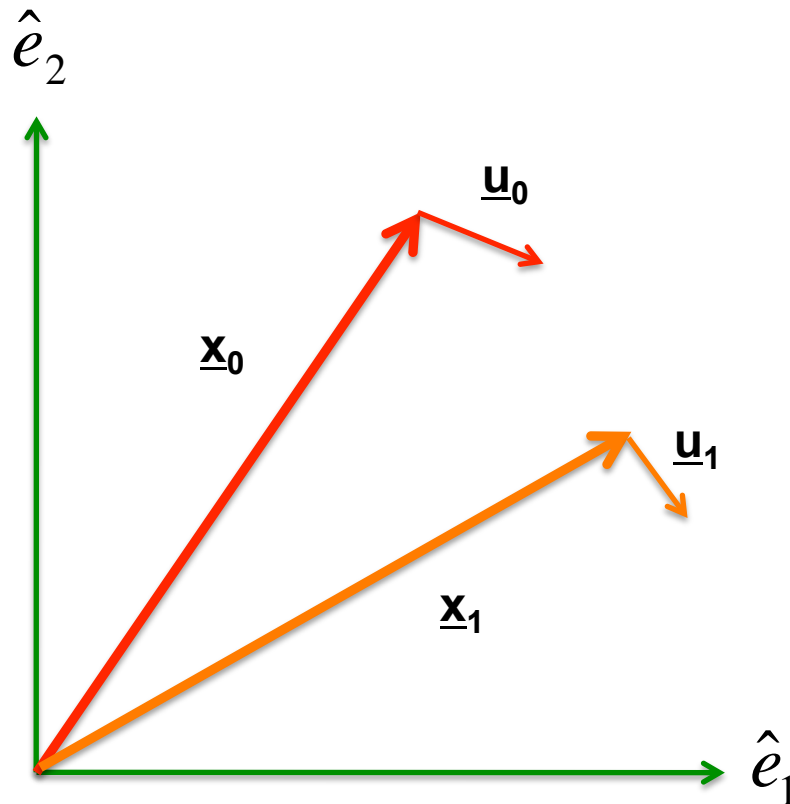
- displacement = rigid body motion + deformation
- displacement = translation + rotation + deformation

# Motion in General



- displacement = rigid body motion + deformation
- displacement = translation + rotation + deformation

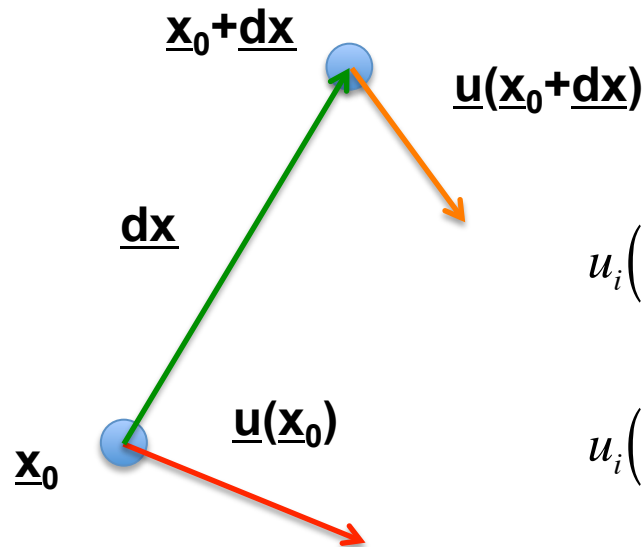
# Displacement



- Every point has some displacement
- $\underline{u}(x_1, x_2, x_3) = \underline{u}(\underline{x})$
- How do we differentiate between rigid motion and deformation?
  - Consider the motion of neighboring points

# Displacement

- Use a Taylor Series expansion to relate the two displacements:



$$u_i(\underline{x}_0 + \underline{dx}) = u_i(\underline{x}_0) + \left( \frac{\partial u_i}{\partial x_j} \right)_{x=\underline{x}_0} dx_j + \dots$$

$$u_i(\underline{x}_0 + \underline{dx}) = u_i(\underline{x}_0) + \left( \frac{\partial u_i}{\partial x_1} \right) dx_1 + \left( \frac{\partial u_i}{\partial x_2} \right) dx_2 + \left( \frac{\partial u_i}{\partial x_3} \right) dx_3$$

This is a set of 3 equations, for  $i = 1, 2, 3$

First term: translation

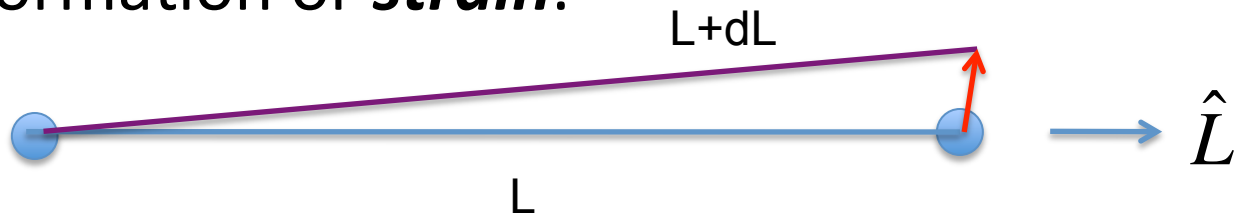
Remaining terms: rotation + strain

There are 9 values  $(\partial u_i / \partial x_j)$  ( $i=1,3; j=1,3$ )



# First a More Physical Point of View

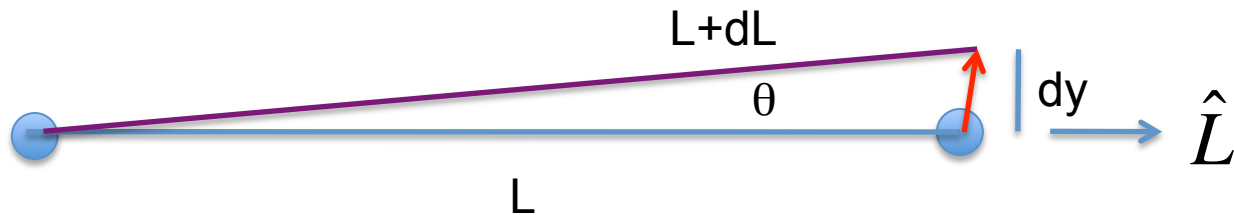
- Let's first look at strain from the point of view of line length. We already know that rotations don't change line lengths, so any line length change comes from deformation or **strain**.



- If one point is displaced relative to another, the line length between them changes from  $L$  to  $L + dL$ 
  - Assume  $dL \ll L$  (for tectonics:  $L \sim \text{km}$ ,  $dL \sim \text{cm}$ )
- Define the strain in the direction of the line to be the fractional change in length, or strain  $\epsilon_{LL} = (dL/L)$

# Shear Strains

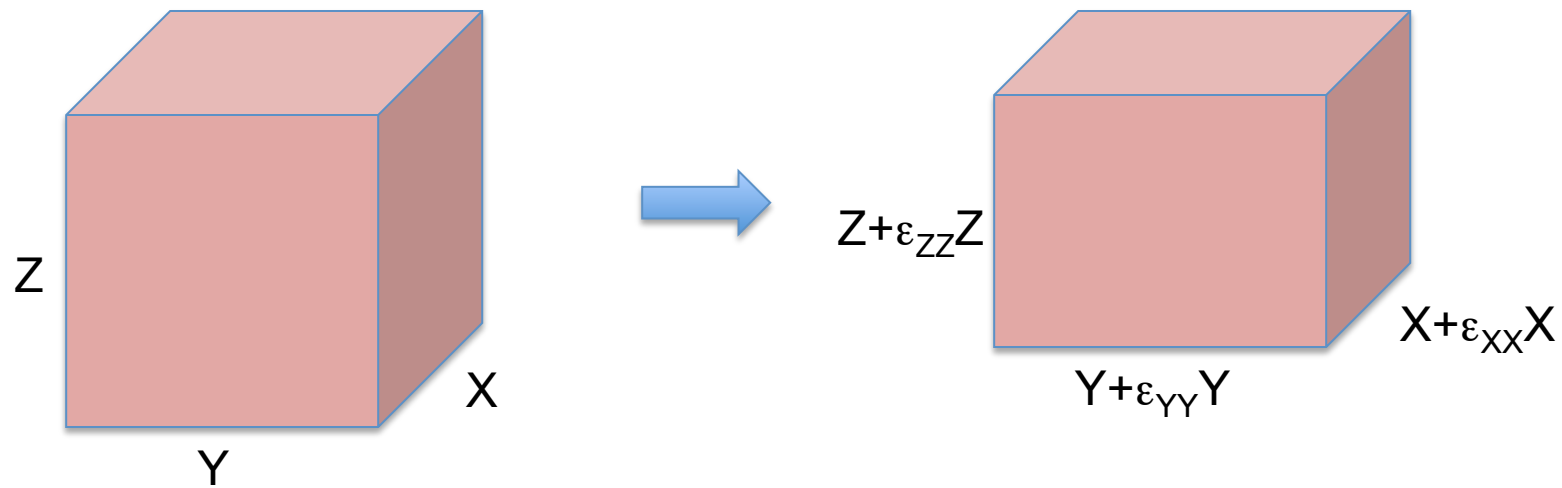
- Strains have direction. In the sketch there may also be a strain in the direction orthogonal to the line, but from this sketch we can't tell if it is deformation or rotation. If we assume no rotation,



- Define the ***shear strain*** to be the displacement in the orthogonal direction divided by the length, or strain  $\epsilon_{\text{shear}} = (dy/L) = \tan\theta \approx \theta$
- There is a better definition of the shear strain, which we will come to in a little bit.

# Dilatation

- Imagine a box that is compressed or extended equally in all dimensions (in the figure, the  $\epsilon$  are negative):



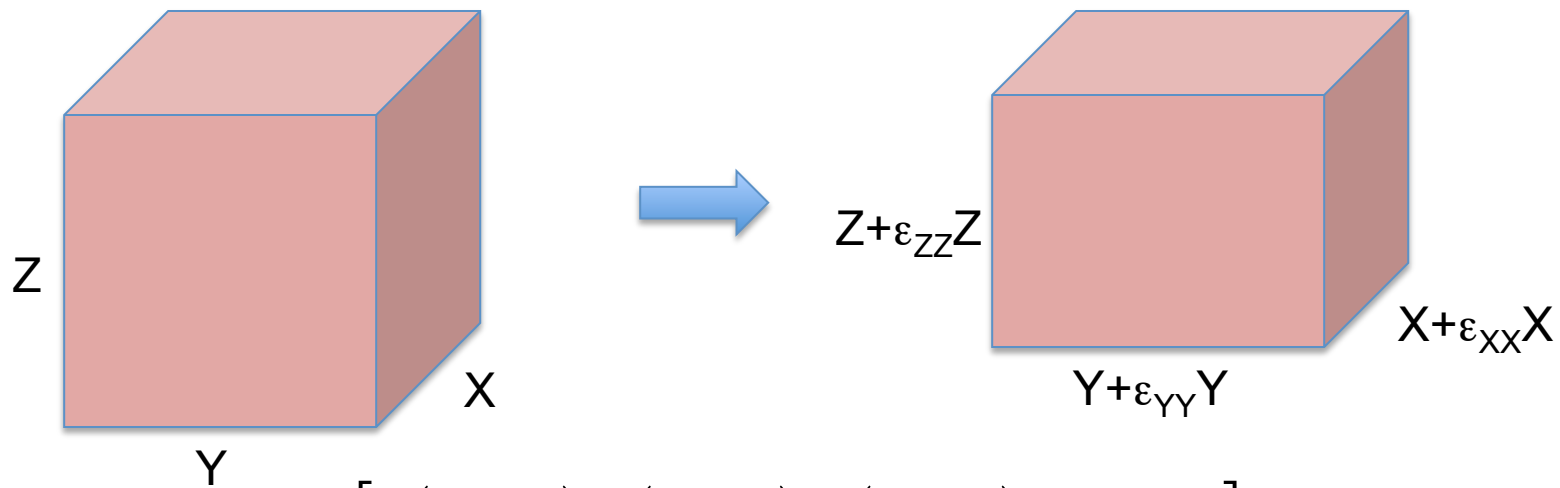
- The length changes (***axial strains***) in the x,y,z directions are
- X:  $dL/L = \epsilon_{XX}$
- Y:  $dL/L = \epsilon_{YY}$
- Z:  $dL/L = \epsilon_{ZZ}$

*Note: positive = extension*

*Geologists often use positive = contraction*

# Dilatation

- Define the dilatation ( $\Delta$ ) as the fractional volume change



$$\Delta = \frac{[X(1 + \epsilon_{XX}) \cdot Y(1 + \epsilon_{YY}) \cdot Z(1 + \epsilon_{ZZ}) - X \cdot Y \cdot Z]}{X \cdot Y \cdot Z}$$

$$\Delta = [(1 + \epsilon_{XX})(1 + \epsilon_{YY})(1 + \epsilon_{ZZ}) - 1]$$

$$\Delta = (1 + \epsilon_{XX} + \epsilon_{YY} + \cancel{\epsilon_{XX}\epsilon_{YY}})(1 + \epsilon_{ZZ}) - 1$$

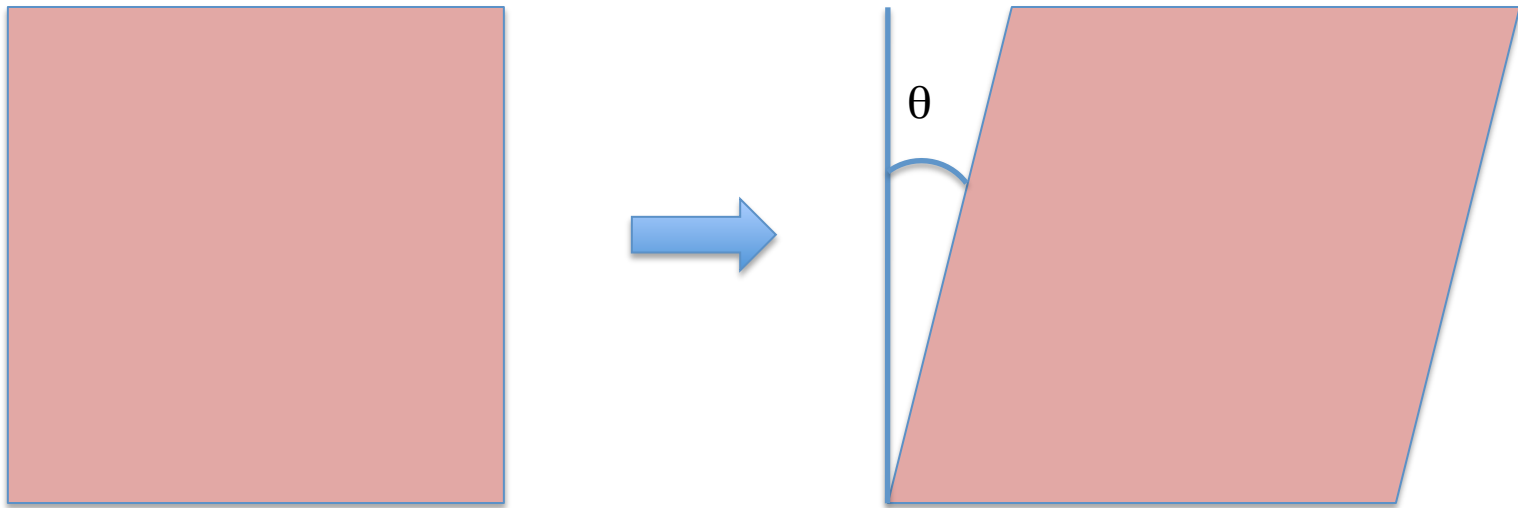
$$\Delta = 1 + \epsilon_{XX} + \epsilon_{YY} + \epsilon_{ZZ} - 1$$

$$\Delta = \epsilon_{XX} + \epsilon_{YY} + \epsilon_{ZZ}$$

*Note: seismic P waves are traveling oscillations of  $\Delta$*

# Shear Strain

- Shear deformation is also very common. Define the shear strain in terms of the angle change of a side:

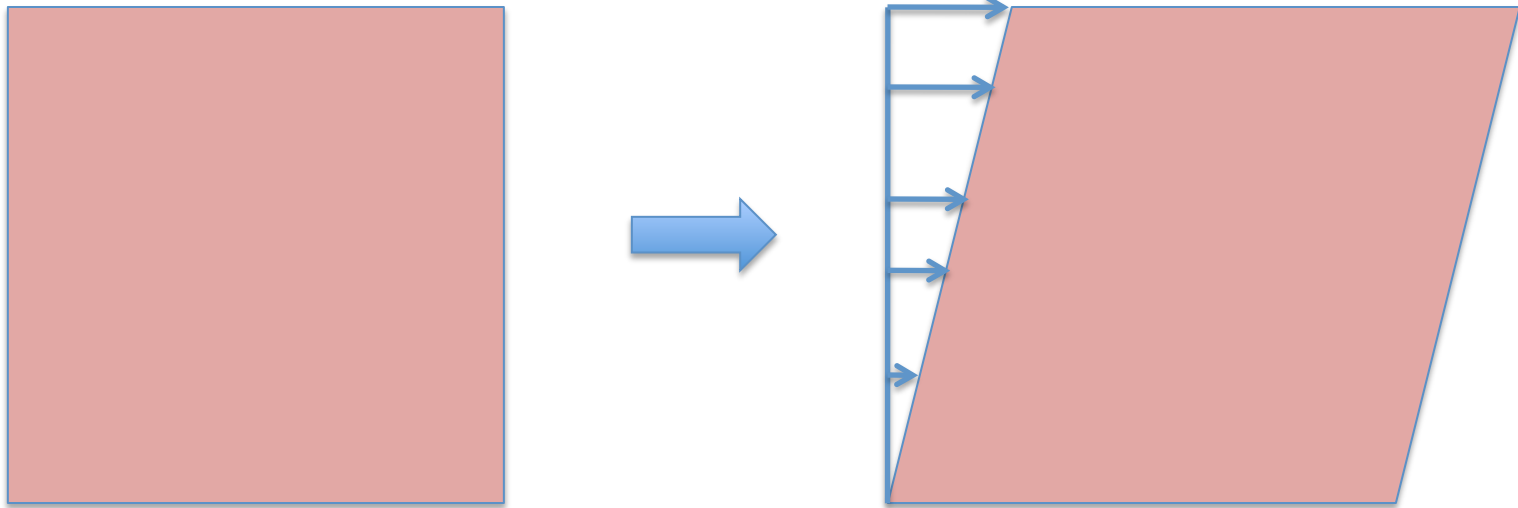


- Shear strain =  $\theta$

*simple shear*

# Shear Strain

- Shear deformation is also very common.  
Define the shear strain in terms of the angle change of a side:

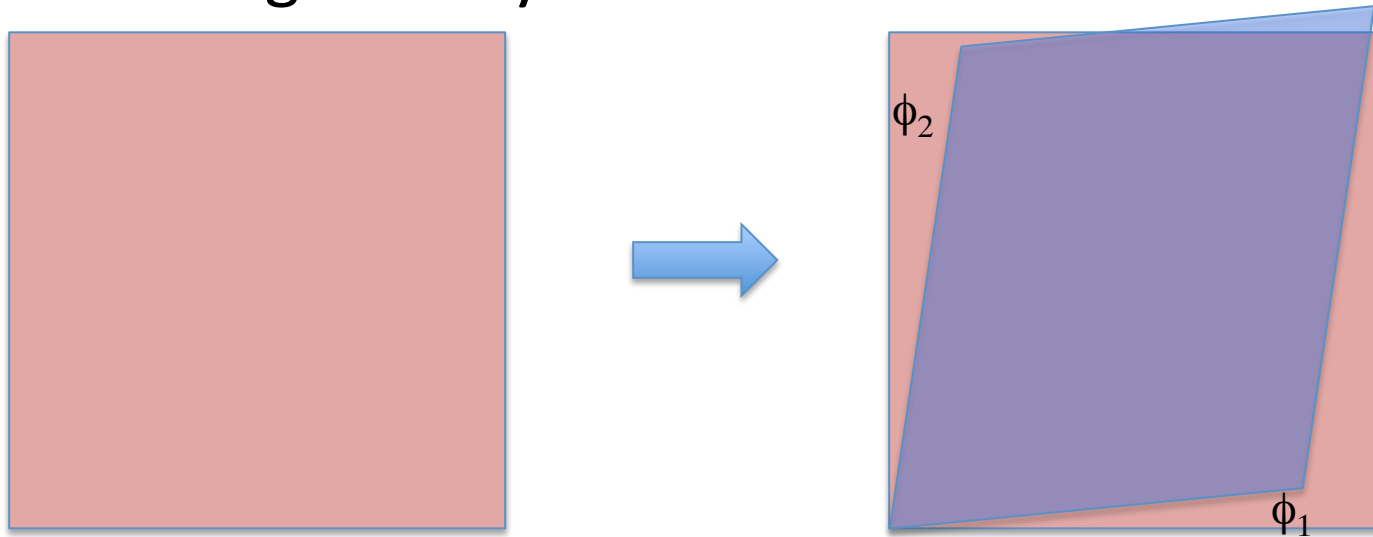


- Shear strain =  $\theta$

*simple shear*

# Shear Strain

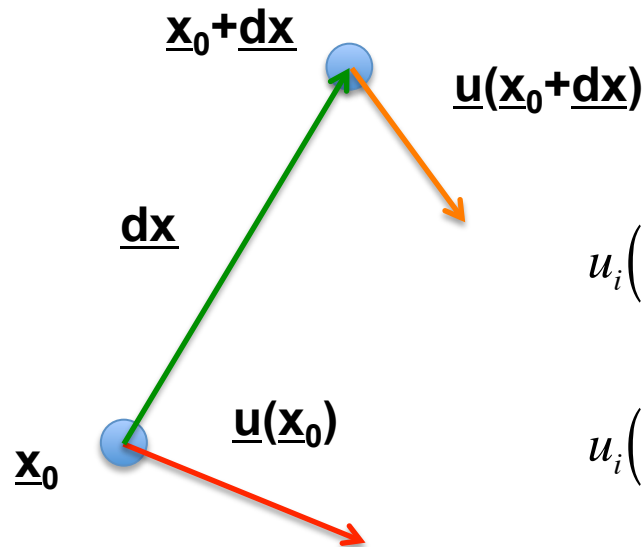
- A bit more generally:



- Shear strain =  $\phi_1 + \phi_2$
- If  $\phi_1 = \phi_2$ , then the strain is ***pure shear***
- Simple shear and pure shear differ only by a rigid rotation

# Displacement

- Use a Taylor Series expansion to relate the two displacements:



$$u_i(\underline{x}_0 + \underline{dx}) = u_i(\underline{x}_0) + \left( \frac{\partial u_i}{\partial x_j} \right)_{x=\underline{x}_0} dx_j + \dots$$

$$u_i(\underline{x}_0 + \underline{dx}) = u_i(\underline{x}_0) + \left( \frac{\partial u_i}{\partial x_1} \right) dx_1 + \left( \frac{\partial u_i}{\partial x_2} \right) dx_2 + \left( \frac{\partial u_i}{\partial x_3} \right) dx_3$$

This is a set of 3 equations, for  $i = 1, 2, 3$

First term: translation

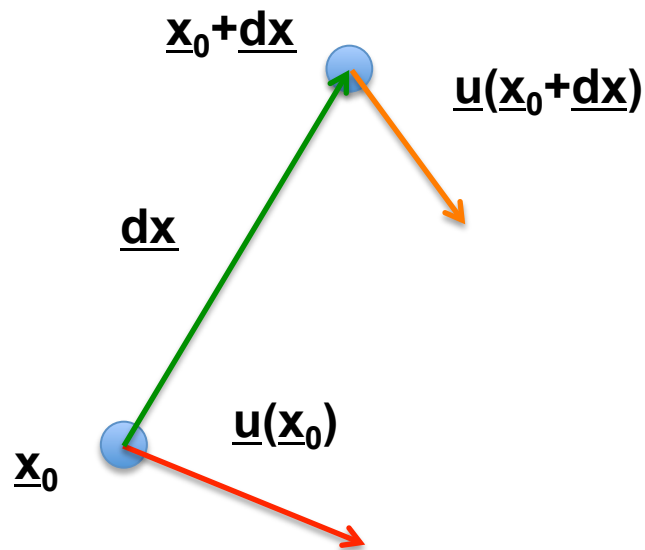
Remaining terms: rotation + strain

There are 9 values  $(\partial u_i / \partial x_j)$  ( $i=1,3; j=1,3$ )



# Displacement Gradient Tensor

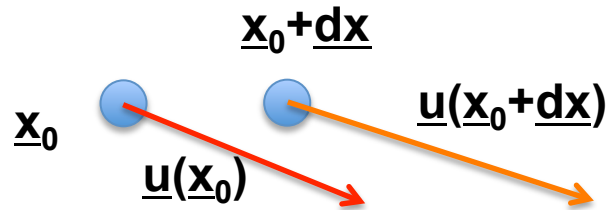
- These 9 values describe the deformation (strain) and rotation together:



$$\begin{bmatrix} \frac{\partial u_1}{\partial x_1} & \frac{\partial u_1}{\partial x_2} & \frac{\partial u_1}{\partial x_3} \\ \frac{\partial u_2}{\partial x_1} & \frac{\partial u_2}{\partial x_2} & \frac{\partial u_2}{\partial x_3} \\ \frac{\partial u_3}{\partial x_1} & \frac{\partial u_3}{\partial x_2} & \frac{\partial u_3}{\partial x_3} \end{bmatrix}$$

# Extension in $\mathbf{e}_1$ direction

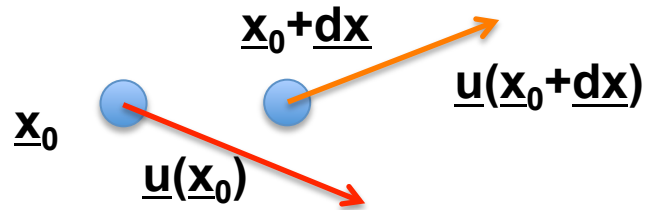
- These 9 values describe the deformation (strain) and rotation together:



$$\begin{bmatrix} \frac{\partial u_1}{\partial x_1} & \frac{\partial u_1}{\partial x_2} & \frac{\partial u_1}{\partial x_3} \\ \frac{\partial u_2}{\partial x_1} & \frac{\partial u_2}{\partial x_2} & \frac{\partial u_2}{\partial x_3} \\ \frac{\partial u_3}{\partial x_1} & \frac{\partial u_3}{\partial x_2} & \frac{\partial u_3}{\partial x_3} \end{bmatrix}$$

# Shear in $\mathbf{e}_2$ direction

- These 9 values describe the deformation (strain) and rotation together:



$$\begin{bmatrix} \frac{\partial u_1}{\partial x_1} & \frac{\partial u_1}{\partial x_2} & \frac{\partial u_1}{\partial x_3} \\ \frac{\partial u_2}{\partial x_1} & \frac{\partial u_2}{\partial x_2} & \frac{\partial u_2}{\partial x_3} \\ \frac{\partial u_3}{\partial x_1} & \frac{\partial u_3}{\partial x_2} & \frac{\partial u_3}{\partial x_3} \end{bmatrix}$$

# Separating Rotation and Strain

- We can define this to be the sum of two tensors, a strain tensor and a rotation tensor. The rotation part will be anti-symmetric (remember the rotation matrix), and the strain part will be symmetric:

$$u_i(\underline{x}_0 + \underline{dx}) = u_i(\underline{x}_0) + \frac{1}{2} \left( \frac{\partial u_i}{\partial x_j} + \frac{\partial u_j}{\partial x_i} \right) dx_j + \frac{1}{2} \left( \frac{\partial u_i}{\partial x_j} - \frac{\partial u_j}{\partial x_i} \right) dx_j$$

$$\begin{bmatrix} \frac{\partial u_1}{\partial x_1} & \frac{\partial u_1}{\partial x_2} & \frac{\partial u_1}{\partial x_3} \\ \frac{\partial u_2}{\partial x_1} & \frac{\partial u_2}{\partial x_2} & \frac{\partial u_2}{\partial x_3} \\ \frac{\partial u_3}{\partial x_1} & \frac{\partial u_3}{\partial x_2} & \frac{\partial u_3}{\partial x_3} \end{bmatrix} = \begin{bmatrix} \frac{\partial u_1}{\partial x_1} & \frac{1}{2} \left( \frac{\partial u_1}{\partial x_2} + \frac{\partial u_2}{\partial x_1} \right) & \frac{1}{2} \left( \frac{\partial u_1}{\partial x_3} + \frac{\partial u_3}{\partial x_1} \right) \\ \frac{1}{2} \left( \frac{\partial u_1}{\partial x_2} + \frac{\partial u_2}{\partial x_1} \right) & \frac{\partial u_2}{\partial x_2} & \frac{1}{2} \left( \frac{\partial u_2}{\partial x_3} + \frac{\partial u_3}{\partial x_2} \right) \\ \frac{1}{2} \left( \frac{\partial u_1}{\partial x_3} + \frac{\partial u_3}{\partial x_1} \right) & \frac{1}{2} \left( \frac{\partial u_2}{\partial x_3} + \frac{\partial u_3}{\partial x_2} \right) & \frac{\partial u_3}{\partial x_3} \end{bmatrix} + \begin{bmatrix} 0 & \frac{1}{2} \left( \frac{\partial u_1}{\partial x_2} - \frac{\partial u_2}{\partial x_1} \right) & \frac{1}{2} \left( \frac{\partial u_1}{\partial x_3} - \frac{\partial u_3}{\partial x_1} \right) \\ \frac{1}{2} \left( \frac{\partial u_2}{\partial x_1} - \frac{\partial u_1}{\partial x_2} \right) & 0 & \frac{1}{2} \left( \frac{\partial u_2}{\partial x_3} - \frac{\partial u_3}{\partial x_2} \right) \\ \frac{1}{2} \left( \frac{\partial u_3}{\partial x_1} - \frac{\partial u_1}{\partial x_3} \right) & \frac{1}{2} \left( \frac{\partial u_3}{\partial x_2} - \frac{\partial u_2}{\partial x_3} \right) & 0 \end{bmatrix}$$

# Strain Tensor

Axial Strains

Shear Strains

The diagram shows the strain tensor matrix with its components categorized into axial and shear strains. The matrix is enclosed in large square brackets. The diagonal elements, representing axial strains, are highlighted in an orange oval. The off-diagonal elements, representing shear strains, are highlighted in a green oval. The matrix is symmetric.

$$\begin{bmatrix} \frac{\partial u_1}{\partial x_1} & \frac{1}{2} \left( \frac{\partial u_1}{\partial x_2} + \frac{\partial u_2}{\partial x_1} \right) & \frac{1}{2} \left( \frac{\partial u_1}{\partial x_3} + \frac{\partial u_3}{\partial x_1} \right) \\ \frac{1}{2} \left( \frac{\partial u_1}{\partial x_2} + \frac{\partial u_2}{\partial x_1} \right) & \frac{\partial u_2}{\partial x_2} & \frac{1}{2} \left( \frac{\partial u_2}{\partial x_3} + \frac{\partial u_3}{\partial x_2} \right) \\ \frac{1}{2} \left( \frac{\partial u_1}{\partial x_3} + \frac{\partial u_3}{\partial x_1} \right) & \frac{1}{2} \left( \frac{\partial u_2}{\partial x_3} + \frac{\partial u_3}{\partial x_2} \right) & \frac{\partial u_3}{\partial x_3} \end{bmatrix}$$

# Rotation Tensor

$$\begin{bmatrix} 0 & \frac{1}{2}\left(\frac{\partial u_1}{\partial x_2} - \frac{\partial u_2}{\partial x_1}\right) & \frac{1}{2}\left(\frac{\partial u_1}{\partial x_3} - \frac{\partial u_3}{\partial x_1}\right) \\ \frac{1}{2}\left(\frac{\partial u_2}{\partial x_1} - \frac{\partial u_1}{\partial x_2}\right) & 0 & \frac{1}{2}\left(\frac{\partial u_2}{\partial x_3} - \frac{\partial u_3}{\partial x_2}\right) \\ \frac{1}{2}\left(\frac{\partial u_3}{\partial x_1} - \frac{\partial u_1}{\partial x_3}\right) & \frac{1}{2}\left(\frac{\partial u_3}{\partial x_2} - \frac{\partial u_2}{\partial x_3}\right) & 0 \end{bmatrix}$$

*In terms of our earlier  
coordinate rotation matrix:*

$$\begin{bmatrix} 0 & \frac{1}{2}\left(\frac{\partial u_1}{\partial x_2} - \frac{\partial u_2}{\partial x_1}\right) & \frac{1}{2}\left(\frac{\partial u_1}{\partial x_3} - \frac{\partial u_3}{\partial x_1}\right) \\ \frac{1}{2}\left(\frac{\partial u_2}{\partial x_1} - \frac{\partial u_1}{\partial x_2}\right) & 0 & \frac{1}{2}\left(\frac{\partial u_2}{\partial x_3} - \frac{\partial u_3}{\partial x_2}\right) \\ \frac{1}{2}\left(\frac{\partial u_3}{\partial x_1} - \frac{\partial u_1}{\partial x_3}\right) & \frac{1}{2}\left(\frac{\partial u_3}{\partial x_2} - \frac{\partial u_2}{\partial x_3}\right) & 0 \end{bmatrix} = (\beta_{ij})_{\underline{x} - \underline{x}}$$